

Above-light-line Nonlinear Surface Polaritons near a Conductive Interface: Threshold Case

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Abstract

We investigate the TM-polarized nonlinear surface polaritons (NLSP) propagating along a guided structure consisting of a magnetic optically linear medium and a non-magnetic optically nonlinear medium with saturable permittivity separated by a flat conductive layer of zero thickness. We consider those values of hosting media bulk material parameters for which the NLSP existence (for zero sheet conductance) has threshold character with respect to the waves intensity. Based on the exact solution of Maxwell's equations we show that the energy and propagation properties of the NLSP near the above-light-line condition ($0 < n < 1$) depend considerably on the surface conductivity of the layer, even the threshold character of the NLSP can be lost; for certain sheet conductance values these waves can exist in a linear limit. The NLSP propagation constant is defined by both the surface conductivity and field intensity and can be varied in a wide diapason, which gives an opportunity to obtain and control the important for quantum information processing $n \approx 0$ condition. For a chosen value of the NLSP propagation constant the NLSP field intensity and energy flux decrease when the surface conductivity grows; saturation of the nonlinear permittivity leads to an increase of the NLSP energy flux compared with Kerr-like nonlinearity. The results are potentially applicable for the development of tunable optical switches of the $n \approx 0$ condition in quantum photonic circuits.

1. Introduction

Surface plasmon polaritons (SPPs) are electromagnetic waves coupled to free electron oscillations traveling, typically, along a metal–dielectric interface [1]. The unique features the waves find promising applications in control and manipulation of light in nanometer scale [1,2]. Subwavelength localization of SPP field around the interface allows to reduce dimensions of plasmonic devices under diffraction limit [3]; strong enhancement of the SPP localized field enlarges nonlinear response of their guiding structure [4]; the SPP field distribution with peak at the interface leads to high sensitivity of the waves to any changes in the interface properties [5]. Surface polaritons were extensively investigated at the surfaces of dielectrics [1,6], ferromagnetics [7], magnetic-semiconductor superlattices [8], metamaterials of hyperbolic [9] and nonlinear [10,11] types, etc. A lot of SPP-based plasmonic

devices and structures have been suggested and fabricated: modulators [12, 13], switches [13], waveguides [14], nanolasers [15], logic gates [16], chemical and biological sensors [17]. Control of the surface polaritons properties is a significant challenge in the creation of plasmonic devices; utilizing graphene in SPP guided structures seems to be the best answer [18]. Graphene is a two-dimensional (2D) material with a single layer of carbon atoms whose high conductivity at infrared frequencies combines with its unique property of being electrically or chemically tunable [19, 20]. Recently proposed graphene plasmonic devices include modulators [21, 22], photodetectors [23], slow-light components [24], filters [25], polarizers [26], and biosensors [27]. Application of surface polaritons for mediating inter-emitter interactions in quantum information processing forms a novel and rapidly growing field of research [28, 29]. It includes realization of quantum interference between single plasmons in a nanoscale circuit [30], strong coupling between surface polaritons and quantum emitters [31] large Purcell enhancement [32], single plasmon generation [33] and detection [34]. Interactions efficiency of quantum emitters can be significantly increased in zero-index plasmonic structures. Stretching of the photon wavelength near the $n \approx 0$ condition enlarges the spatial scale of inter-emitter interactions and opens a way to control the interactions of quantum emitters which are separated at long distances with no need for their precise positioning [35-37].

Properties of surface polaritons are defined by the material parameters of their guiding structures. Recent progress in optical metamaterials confirm that every desirable value of their material parameters can be obtained, including negative and near-zero values of the magnetic permeability [38-47]. Magneto-optical metamaterials are created as substrates doped with metal nanoparticles of different shape and size [41-44] or low-loss all-dielectric structures [45]. Nonlinear metamaterials contain layers of metal and nonlinear media, or nonlinear media doped with metal nanoparticles [48]. The local field enhancement in the metamaterials can enlarge their nonlinearities and lead to saturation of the nonlinear permittivity. Surface polaritons in such artificially designed media obtain unexpected features.

Above-light-line surface polaritons whose phase speed exceeds the speed of light in vacuum can exist near the surface of an epsilon-near-zero (ENZ) metamaterial [49]. The effective refractive index of the polaritons is less than unity, the photon wavelength is larger and the phase advance

is smaller than the same values for light in free space. Such properties lead to homogenization of the polaritons field distribution in the direction of propagation, which can be used to increase the interactions efficiency of quantum emitters in quantum information processing [37]. Contrary to the properties of bulk and guided electromagnetic waves in near-zero-index structures [35-37,50,51], for the surface polaritons in non-magnetic media, there exists a long-wavelength cutoff [52] that is an obstacle in achieving the important for quantum information processing $n \approx 0$ condition. This difficulty can be overcome in magnetic media by a proper selection of the media dielectric permittivity and magnetic permeability values [53].

TM-polarized nonlinear surface polaritons (NLSP) can propagate along the surface of an optically nonlinear medium. Besides TM-polarized NLSP that exist in linear approximation there are ones whose existence becomes possible only in the nonlinear case and that have threshold character with respect to the NLSP field intensity or the NLSP energy flow. Threshold NLSP propagating along the interface of non-magnetic ENZ media obtain above-light-line features [54]. Implementation of conductive layers into waveguide structures can be used to provide tunable optical switches on deep-subwavelength scale. Plasmonic structure containing semi-infinite graphene sheet arrays embedded in the linear dielectric was proposed as a tunable optical switch whose threshold power is controlled by the chemical potential of graphene [55]; plasmonic structure with a single graphene monolayer demonstrates features of tunable optical bistable device with ultralow threshold value [56].

In this work we consider the threshold NLSP propagating along the guided structure consisting of a magnetic optically linear ENZ medium and a non-magnetic defocusing ENZ medium with saturable permittivity separated by a flat conductive layer of zero thickness. Based on the exact solution of Maxwell's equations we investigate the NLSP energy and propagation properties for the above-light-line condition $0 < n < 1$, namely the dependences of the NLSP propagating constant vs. the NLSP total energy flux and the dependences of the NLSP propagating constant vs. the permittivity at the surface of the nonlinear medium for different values of sheet conductance for both saturable and Kerr-like nonlinearity.

2. Statement of the problem

We study the NLSP propagating in x -direction in the following lossless structure:

- semi-infinite optically linear magnetic medium (region $z < 0$; dielectric permittivity ε_1 , magnetic permeability μ_1);
- conductive layer of zero thickness ($z=0$; sheet conductance σ);
- semi-infinite optically nonlinear non-magnetic medium (region $z > 0$; intensity-dependent permittivity ε).

The problem is assumed to be uniform in y -direction. The electric and magnetic NLSP fields are expressed by

$$\vec{E}(\vec{r}, t) = \frac{1}{2} (\vec{E}(z) \exp i(kx - \omega t) + c.c.), \quad \vec{E}(z) = (E_x, 0, E_z), \quad (1)$$

$$\vec{H}(\vec{r}, t) = \frac{1}{2} (\vec{H}(z) \exp i(kx - \omega t) + c.c.), \quad \vec{H}(z) = (0, H, 0). \quad (2)$$

The sheet conductance value of the layer is supposed to be purely imaginary: $\sigma = i\sigma''$ (it corresponds to graphene, at optical frequencies graphene is almost lossless and its conductivity can be considered a purely imaginary value [57]). The function which models the permittivity of the nonlinear medium is as follows:

$$\varepsilon = \varepsilon_{xx} = \varepsilon_{zz} = \varepsilon_0 + \varepsilon_2 (|\vec{E}|^2) = \varepsilon_0 + \frac{a(|E_x|^2 + |E_z|^2)}{1 + b(|E_x|^2 + |E_z|^2)}, \quad (3)$$

here ε_0 and $\varepsilon_2 (|\vec{E}|^2)$ are the linear and the nonlinear parts of the permittivity. The quantity b is supposed to be non-negative; the signs of ε_0 , a , ε_1 , and μ_1 are arbitrary. The saturation level of the intensity-dependent permittivity is $\varepsilon_s = \varepsilon_0 + a/b$. The permittivity is Kerr-like for weak fields ($b|\vec{E}|^2 \ll 1$) and saturates for strong fields.

3. Theory

Maxwell's equations for the NLSP field components give:

$$\frac{dH}{d\zeta} = i\varepsilon E_x \quad (4a)$$

$$nH = -\varepsilon E_z \quad (4b)$$

$$\frac{dE_x}{d\zeta} - inE_z = i\mu H, \quad (4c)$$

where $n \equiv k/q$ is the dimensionless propagation constant and $\zeta \equiv zq$ is the dimensionless coordinate; $q = \omega/c$ is the wavevector of light in free space.

The solution of the system of Eq.(4a) to Eq.(4c) for the linear magnetic medium is as follows:

$$H(\zeta) = H_L(0) \exp(k_1 \zeta), \quad (5)$$

$$k_1 = +(n^2 - \varepsilon_1 \mu_1)^{1/2},$$

here $H_L(0)$ is the magnitude of the NLSP magnetic field at the surface of the linear medium.

The solution of the system of Eq.(4a) to Eq.(4c) for the nonlinear medium (equations (2.10) and (2.13b) of

[52]) gives the dependence $H = H(\zeta)$ by expressing the NLSP magnetic field H in terms of dielectric permittivity ε , and the dimensionless coordinate ζ as an integral between the limits $\tilde{\varepsilon}$ and ε , where $\tilde{\varepsilon}$ is the dielectric permittivity at the surface of the nonlinear medium. The evident form of the relationship depends on $I(\varepsilon - \varepsilon_0)$ which is the inverse function of $\varepsilon_2(|\vec{E}|^2)$. Explicit forms of $I(\varepsilon - \varepsilon_0)$ and its integral $J(\varepsilon - \varepsilon_0)$ for the permittivity expressed by equation (3) are defined by Eqs. (9) and (10) from [58]:

$$I(\varepsilon - \varepsilon_0) = (\varepsilon - \varepsilon_0)(a - b(\varepsilon - \varepsilon_0))^{-1}, \quad (6)$$

$$J(\varepsilon - \varepsilon_0) = -\frac{1}{b} \left\{ \varepsilon - \varepsilon_0 + \left[\frac{a}{b} \right] \ln \left[1 - \frac{b}{a} (\varepsilon - \varepsilon_0) \right] \right\}. \quad (7)$$

For the case under consideration the tangential component of the NLSP electric field E_x is continuous across the interface $z=0$, but the NLSP magnetic field has a gap whose value is proportional to the sheet conductance of the layer:

$$H_{NL}(0) - H_L(0) = \sigma \frac{4\pi}{c} E_x(0), \quad (8)$$

where $H_{NL}(0)$ and $E_x(0)$ are the magnitudes of the NLSP magnetic and electric field components at the surface of the nonlinear medium.

Making use of the boundary condition one can obtain the equation for the NLSP propagation constant n :

$$\begin{aligned} & k_1^4 \left[\tilde{\varepsilon} (\tilde{\varepsilon} I(\tilde{\varepsilon} - \varepsilon_0) - J(\tilde{\varepsilon} - \varepsilon_0)) \frac{1}{\varepsilon_1} - I(\tilde{\varepsilon} - \varepsilon_0) \frac{2}{\varepsilon_1^2} \right] \left(\frac{4\pi\sigma''}{c} \right)^2 - \\ & k_1^3 \left[\tilde{\varepsilon} (\tilde{\varepsilon} I(\tilde{\varepsilon} - \varepsilon_0) - J(\tilde{\varepsilon} - \varepsilon_0)) \frac{2}{\varepsilon_1} - I(\tilde{\varepsilon} - \varepsilon_0) \frac{4}{\varepsilon_1^2} \right] \left(\frac{4\pi\sigma''}{c} \right) + \\ & k_1^2 \left[\tilde{\varepsilon} (\tilde{\varepsilon} I(\tilde{\varepsilon} - \varepsilon_0) - J(\tilde{\varepsilon} - \varepsilon_0)) \left\{ 1 + \left(\frac{4\pi\sigma''}{c} \right)^2 \frac{\mu_1}{\varepsilon_1} + \frac{1}{\varepsilon_1^2} \right\} - \right. \\ & \left. I(\tilde{\varepsilon} - \varepsilon_0) \left\{ 2 + \left(\frac{4\pi\sigma''}{c} \right)^2 \left(\frac{2\mu_1}{\varepsilon_1} - \frac{\tilde{\varepsilon}}{\varepsilon_1^2} \right) \right\} \right] - \\ & k_1 \left[\tilde{\varepsilon} (\tilde{\varepsilon} I(\tilde{\varepsilon} - \varepsilon_0) - J(\tilde{\varepsilon} - \varepsilon_0)) 2 \left(\frac{4\pi\sigma''}{c} \right) \mu_1 - \right. \\ & \left. I(\tilde{\varepsilon} - \varepsilon_0) \left\{ 4\mu_1 - 2 \frac{\tilde{\varepsilon}}{\varepsilon_1} \right\} \left(\frac{4\pi\sigma''}{c} \right) \right] + \end{aligned} \quad (9)$$

$$\tilde{\varepsilon} (\tilde{\varepsilon} I(\tilde{\varepsilon} - \varepsilon_0) - J(\tilde{\varepsilon} - \varepsilon_0)) \varepsilon_1 \mu_1 - I(\tilde{\varepsilon} - \varepsilon_0) (2\varepsilon_1 \mu_1 - \tilde{\varepsilon}) = 0.$$

Roots of this four-order equation with respect to the positive $k_1 = +(n^2 - \varepsilon_1 \mu_1)^{1/2}$ give the NLSP propagation constant as a function of both the dielectric permittivity at the surface of the nonlinear medium $\tilde{\varepsilon}$ and the sheet conductance value of the layer σ'' .

For a zero-conductivity case eq. (9) transforms into the following expression:

$$n^2 = \frac{\varepsilon_1 \tilde{\varepsilon}^2 ((\mu_1 \tilde{\varepsilon} - \varepsilon_1) I(\tilde{\varepsilon} - \varepsilon_0) - \mu_1 J(\tilde{\varepsilon} - \varepsilon_0))}{\tilde{\varepsilon} (\tilde{\varepsilon}^2 - \varepsilon_1^2) I(\tilde{\varepsilon} - \varepsilon_0) - (\tilde{\varepsilon}^2 + \varepsilon_1^2) J(\tilde{\varepsilon} - \varepsilon_0)}. \quad (10)$$

The requirement for the NLSP field to vanish for $z \rightarrow +\infty$ and $z \rightarrow -\infty$ gives for polaritons in non-magnetic media a long-wavelength cutoff for the NLSP propagation constant $n^2 > \max(0, \varepsilon_0, \varepsilon_1)$ which is valid in both linear and nonlinear cases [52]. In the case under consideration the cutoff obtains the following form:

$$n^2 > \max(0, \varepsilon_0, \varepsilon_1 \mu_1). \quad (11)$$

For both negative ε_0 and $\varepsilon_1 \mu_1$ the cutoff obtains zero value, and the near-zero index condition $n \approx 0$ can be achieved.

The NLSP total energy flux per unit length in y -direction Q consists of the energy flux in the linear medium Q_L and the energy flux in the nonlinear medium Q_{NL} . The results of [52] allow expressing both these fluxes in terms of $\tilde{\varepsilon}$:

$$Q_L = \frac{cn}{16\pi\varepsilon_1(n^2 - \mu_1\varepsilon_1)^{1/2}q} \frac{\tilde{\varepsilon}}{(2n^2 - \tilde{\varepsilon})} (\tilde{\varepsilon} I(\tilde{\varepsilon} - \varepsilon_0) - J(\tilde{\varepsilon} - \varepsilon_0)) \quad (12)$$

$$Q_{NL} = \text{sgn}(\varepsilon'_\zeta) \frac{c}{8\pi} \frac{n}{q} \int_{\tilde{\varepsilon}}^{\varepsilon_0} d\varepsilon \frac{\varepsilon I(\varepsilon - \varepsilon_0) - J(\varepsilon - \varepsilon_0)}{(2n^2 - \varepsilon) [-2V(\varepsilon)]^{1/2}}, \quad (13)$$

where

$$V(\varepsilon) = -\frac{1}{2} \left[\frac{dH(\varepsilon)}{d\varepsilon} \right]^{-2} \left[\varepsilon^2 I(\varepsilon - \varepsilon_0) - n^2 H^2(\varepsilon) \right], \quad (14)$$

and the expression for the NLSP magnetic field is as follows:

$$H^2(\varepsilon) = \frac{\varepsilon}{2n^2 - \varepsilon} (\varepsilon I(\varepsilon - \varepsilon_0) - J(\varepsilon - \varepsilon_0)). \quad (15)$$

4. Specific examples

The calculations of the NLSP properties were carried out for the following values of parameters:

- free-space wavelength: $\lambda = 2\pi c / \omega = 1.06 \mu\text{m}$;
- the linear medium: permittivity $\epsilon_1 = 0.081$, permeability $\mu_1 = -0.6$;
- the nonlinear medium: linear part of permittivity $\epsilon_0 = -0.08$, Kerr coefficient $a = -4.4 \times 10^{-20} \text{ m}^2 / \text{V}^2$ (that is a typical value of nonlinear optical materials). In linear approximation ($a = 0$) there are no surface polaritons for such material parameters (the squared linear propagation constant $n_L^2 = (\epsilon_0 \epsilon_1 (\mu_1 \epsilon_0 - \epsilon_1)) / (\epsilon_0^2 - \epsilon_1^2)$ [59] is negative), but their existence becomes possible in the nonlinear case and the corresponding NLSP has threshold behavior with respect to the NLSP energy flux and to the permittivity value at the surface of the nonlinear medium.

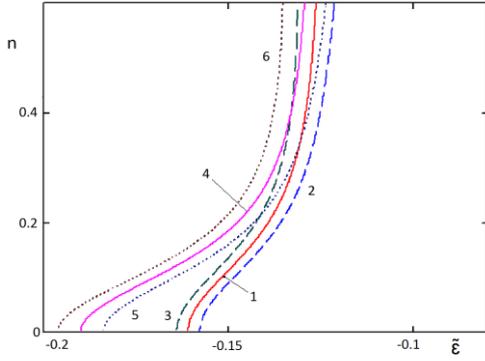


Figure 1. The NLSP propagation constant n vs. the permittivity at the surface of the nonlinear medium $\tilde{\epsilon}$. The saturation level is $\epsilon_s = -0.2$: (1) $\sigma'' = 0$, (2) and (3) $\sigma'' = -10^{-5} \text{ S}$; the saturation level is $\epsilon_s = -\infty$: (4) $\sigma'' = 0$, (5) and (6) $\sigma'' = -10^{-5} \text{ S}$.

Figure 1 presents the dependences of the NLSP propagation constant vs. the permittivity at the surface of the nonlinear medium $\tilde{\epsilon}$ for the saturation levels $\epsilon_s = -0.2$ and $\epsilon_s = -\infty$ for the sheet conductance values $\sigma'' = 0$ and $\sigma'' = -10^{-5} \text{ S}$ (a typical value for graphene at infrared frequencies [57]). For the case of zero conductivity the propagation constant values are defined by eq. (10) and there exists a single trace ((1) for $\epsilon_s = -0.2$ and (4) for $\epsilon_s = -\infty$) which originates from the $n = 0$ and approaches infinity. The traces do not intersect the axis n , so such NLSP do not exist in linear limit. For the case $\sigma'' \neq 0$ the propagation constant values are defined by eq.(9) and the abovementioned zero-conductivity traces are surrounded by the traces for $\sigma'' = -10^{-5} \text{ S}$: ((2) and (3) for $\epsilon_s = -0.2$; (5) and (6) for $\epsilon_s = -\infty$). The traces (2) and (5) tend to a single

point at the axis n , which gives the NLSP propagation constant value in the linear limit.

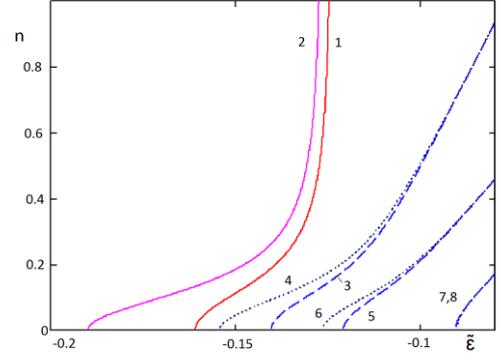


Figure 2. The NLSP propagation constant n vs. the permittivity at the surface of the nonlinear medium $\tilde{\epsilon}$. For the saturation level $\epsilon_s = -0.2$: (1) $\sigma'' = 0$, (3) $\sigma'' = -6 \cdot 10^{-5} \text{ S}$, (5) $\sigma'' = -11 \cdot 10^{-5} \text{ S}$, (7) $\sigma'' = -18 \cdot 10^{-5} \text{ S}$; for the saturation level $\epsilon_s = -\infty$: (2) $\sigma'' = 0$, (4) $\sigma'' = -6 \cdot 10^{-5} \text{ S}$, (6) $\sigma'' = -11 \cdot 10^{-5} \text{ S}$, (8) $\sigma'' = -18 \cdot 10^{-5} \text{ S}$.

Figure 2 shows that with the increase of conductivity the corresponding intersection points approach smaller values of the propagation constant: traces (3) for $\epsilon_s = -0.2$ and (4) for $\epsilon_s = -\infty$ correspond to $\sigma'' = -6 \cdot 10^{-5} \text{ S}$; traces (5) for $\epsilon_s = -0.2$ and (6) for $\epsilon_s = -\infty$ correspond to $\sigma'' = -11 \cdot 10^{-5} \text{ S}$; traces (7) for $\epsilon_s = -0.2$ and (8) for $\epsilon_s = -\infty$ correspond to $\sigma'' = -18 \cdot 10^{-5} \text{ S}$. Figures 1 and 2 illustrate that NLSP properties depend considerably on the surface conductivity of the layer; for a certain value of sheet conductance the waves can exist even for near-zero field intensity.

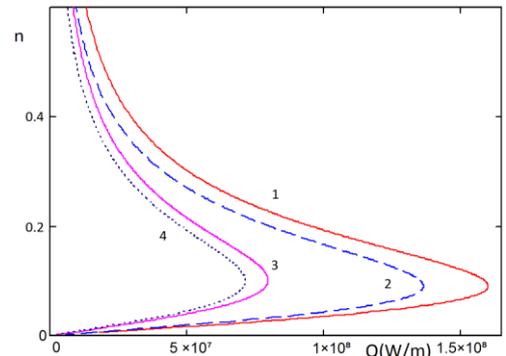


Figure 3. The NLSP propagation constant n vs. the NLSP total energy flux Q . The saturation level is $\epsilon_s = -0.2$: (1)

$\sigma''=0$, (2) $\sigma''=-10^{-5}$ S ; the saturation level is $\varepsilon_s=-\infty$: (3) $\sigma''=0$, (4) $\sigma''=-10^{-5}$ S.

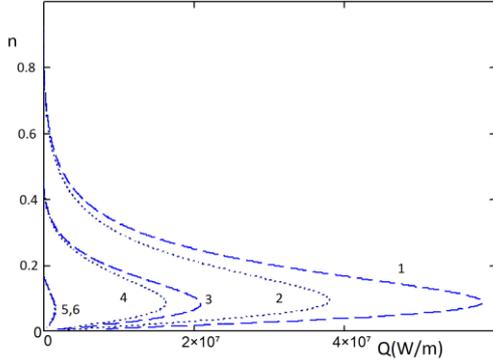


Figure 4. The NLSP propagation constant n vs. the NLSP total energy flux Q . The saturation level is $\varepsilon_s = -0.2$: (1) $\sigma'' = -6 \cdot 10^{-5}$ S, (3) $\sigma'' = -11 \cdot 10^{-5}$ S, (5) $\sigma'' = -18 \cdot 10^{-5}$ S; the saturation level is $\varepsilon_s = -\infty$: (2) $\sigma'' = -6 \cdot 10^{-5}$ S, (4) $\sigma'' = -11 \cdot 10^{-5}$ S, (6) $\sigma'' = -18 \cdot 10^{-5}$ S.

The dependences of the NLSP propagation constant vs. the NLSP energy flux $n = n(Q)$ are presented in Fig.3. The zero conductivity traces (1) for $\varepsilon_s = -0.2$ and (3) for $\varepsilon_s = -\infty$ correspond to the traces (1) and (4) from Figure 1. The traces for non-zero conductivity (2) for $\varepsilon_s = -0.2$ and (4) for $\varepsilon_s = -\infty$ correspond to traces (2) and (5) from Fig. 1. All the traces originate from $n = 0$; the zero conductivity ones approach infinity, the non-zero conductivity ones intersect the axis n . Figure 4 presents the dependences $n = n(Q)$ for increased surface conductivity. The growth of the surface conductivity leads to a decrease of the NLSP energy flux; the saturation of the nonlinear permittivity leads to growth of the NLSP energy flux compared with Kerr-like case.

The NLSP combine above-light-line features with threshold character of their excitation and are potentially applicable for the development of optical switch with tunable threshold level for application in quantum information processing circuits.

5. Conclusion

In this work we investigate the threshold NLSP propagating along the guided structure consisting of a magnetic optically linear ENZ medium and a non-magnetic optically nonlinear (defocusing) ENZ medium with saturable permittivity separated by a flat conductive layer of zero thickness. Based on the exact solution of Maxwell's equations we show that near the above-light-line condition ($0 < n < 1$) the NLSP energy and propagation properties depend considerably on

the sheet conductance value of the layer, even the threshold character of the NLSP can be lost and the NLSP can exist for vanishing field magnitude. The NLSP properties (and even their existence) can be controlled in two independent ways: by field intensity and by surface conductivity. The NLSP propagation constant can be varied in a wide diapason, from infinity to zero in the case of zero conductivity and from the certain finite value to zero in the case of non-zero conductivity, which gives an opportunity to obtain and control the important for quantum information processing $n \approx 0$ condition. For a chosen value of the NLSP propagation constant the NLSP field intensity and energy flux decrease when the surface conductivity grows; saturation of the nonlinear permittivity leads to an increase of the NLSP energy flux compared with Kerr-like nonlinearity.

The results would provide potential application for tunable optical switching of the $n \approx 0$ condition in quantum information processing circuits.

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